

## The TV-ICE model

### OBSERVATION MODEL

Let  $u : \Omega \rightarrow \mathbb{R}_+$  denote an (unobserved) intensity image defined on a discrete domain  $\Omega$ , and  $v$  a random photon-count observation of the ideal image  $u$ , following the **Poisson probability density function**

$$p(v|u) = \prod_{x \in \Omega} \frac{u(x)^{v(x)}}{v(x)!} e^{-u(x)} \propto e^{-\langle u-v \log u, \mathbb{1}_\Omega \rangle},$$

where  $\mathbb{1}_\Omega$  denotes the constant image equal to 1 on  $\Omega$  and  $\langle \cdot, \cdot \rangle$  is the usual inner product on  $\mathbb{R}^\Omega$ .

### MAXIMUM A POSTERIORI (MAP)

The equivalent of the Rudin-Osher-Fatemi (ROF) model [1] in the case of a Poisson noise model corresponds to the unique minimizer of the convex energy  $E(u) = \langle u - v \log u, \mathbb{1}_\Omega \rangle + \lambda \text{TV}(u)$  which is usually reinterpreted from a Bayesian point of view as the unique **maximizer of the posterior density**

$$\pi(u) = p(u|v) \propto p(v|u)p(u)$$

taking  $p(u) \propto e^{-\lambda \text{TV}(u)}$ . The main drawback of this approach is that the restored image generally suffers from the **staircasing effect** due to the non-differentiability of the Total Variation (TV) term.

### LEAST SQUARE ERROR (TV-LSE)

In the case of Gaussian noise, this undesirable staircasing effect can be avoided by considering the **posterior mean** instead of the maximizer of  $p(u|v)$ , that is,

$$\hat{u}_{\text{LSE}} = \mathbb{E}_{u \sim \pi}(u) = \int_{\mathbb{R}^\Omega} u \pi(u) du,$$

which is the image that **reaches the Least Square Error** under  $\pi$  (see [2]). As in the Gaussian case, the numerical computation of  $\hat{u}_{\text{LSE}}$  could be done using a Markov Chain Monte Carlo Metropolis-Hasting algorithm, but would exhibit a **slow convergence rate**.

### ITERATED CONDITIONAL MEANS (TV-ICE)

To overcome this computational limitation, in [3] a new variant was proposed based on the iteration of **conditional marginal posterior means** (in the case of a Gaussian noise). The  $\hat{u}_{\text{ICE}}$  estimate is defined as the limit of the iterative scheme

$$\forall x \in \Omega, u^{n+1}(x) = \mathbb{E}_{u \sim \pi}(u(x) | u(x^c) = u^n(x^c)),$$

noting  $x^c = \Omega \setminus \{x\}$ . Like in the Gaussian case, the iterates are relatively **easy to compute** as they only involve the computation of integrals over  $\mathbb{R}$ .

In order to obtain explicit formulae we need to consider the **anisotropic** version of the discrete total variation

$$\text{TV}(u) = \frac{1}{2} \sum_{x \in \Omega} \sum_{y \in \mathcal{N}_x} |u(x) - u(y)|,$$

where  $\mathcal{N}_x$  represents the 4-neighborhood of the pixel  $x$ . Noting  $a = (a_1, a_2, a_3, a_4)$  the values of  $u^n(\mathcal{N}_x)$  sorted in nondecreasing order (set  $a_0 = 0 \leq a_1 \leq a_2 \leq a_3 \leq a_4 \leq a_5 = +\infty$ ). We obtain

$$u^{n+1}(x) = \frac{\sum_{1 \leq k \leq 5} c_k I_{a_{k-1}, a_k}^{\mu_k, v(x)+1}}{\sum_{1 \leq k \leq 5} c_k I_{a_{k-1}, a_k}^{\mu_k, v(x)}}, \quad (1)$$

where  $c_k, \mu_k$  are simple constants that explicitly depend on  $k, a$  and  $\lambda$ , while  $I_{x,y}^{\mu,p}$  denotes the integral

$$I_{x,y}^{\mu,p} = \int_x^y s^p e^{-\mu s} ds,$$

whose numerical evaluation is closely related to that of lower and upper **incomplete gamma functions**.

## The Poisson TV-ICE algorithm

### NUMERICAL COMPUTATION OF THE POISSON TV-ICE

The practical computation of the TV-ICE recursion raises several **numerical issues**:

- For some values of the parameters, the numerator and the denominator in (1) **cannot be represented by the machine arithmetic** although the **actual value of the ratio is representable**. To solve that issue we represent each integral  $I_{x,y}^{\mu,p}$  under the form  $\rho \times e^\sigma$ .
- An integral  $I_{x,y}^{\mu,p}$  can be computed as the difference between generalized lower ( $\gamma_\mu$ ) or/and upper ( $\Gamma_\mu$ ) incomplete gamma functions:
  - $I_{x,y}^{\mu,p} = \gamma_\mu(p+1, y) - \gamma_\mu(p+1, x)$ ,
  - $I_{x,y}^{\mu,p} = \Gamma_\mu(p+1, x) - \Gamma_\mu(p+1, y)$ ,
  - $I_{x,y}^{\mu,p} = \frac{p!}{\mu^{p+1}} - \gamma_\mu(p+1, x) - \Gamma_\mu(p+1, y)$ .

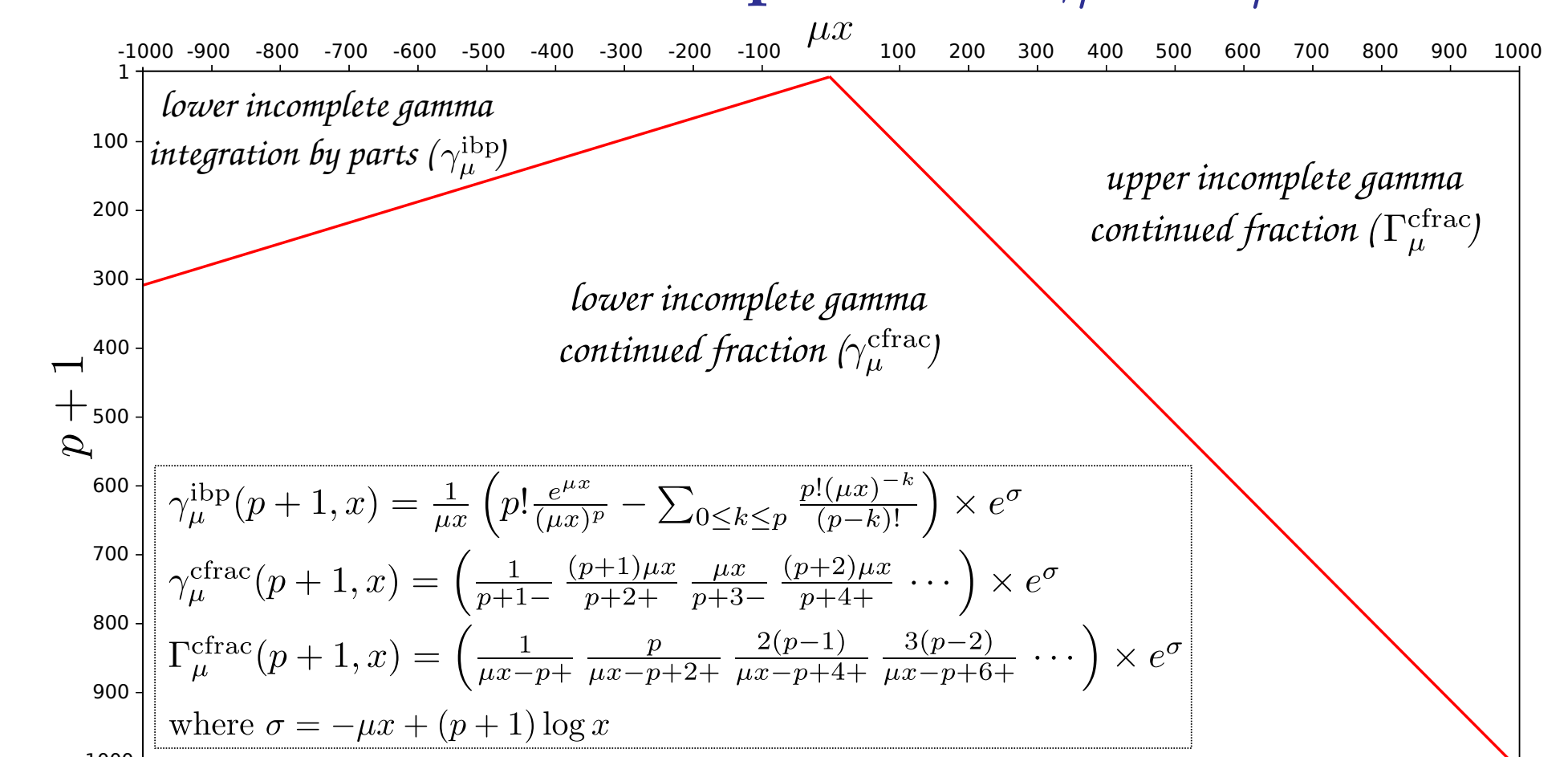
We derived a partition of the plane  $(\mu, p)$  that allows an **efficient computation** of  $\gamma_\mu$  or  $\Gamma_\mu$ .

### Generalized incomplete gamma functions

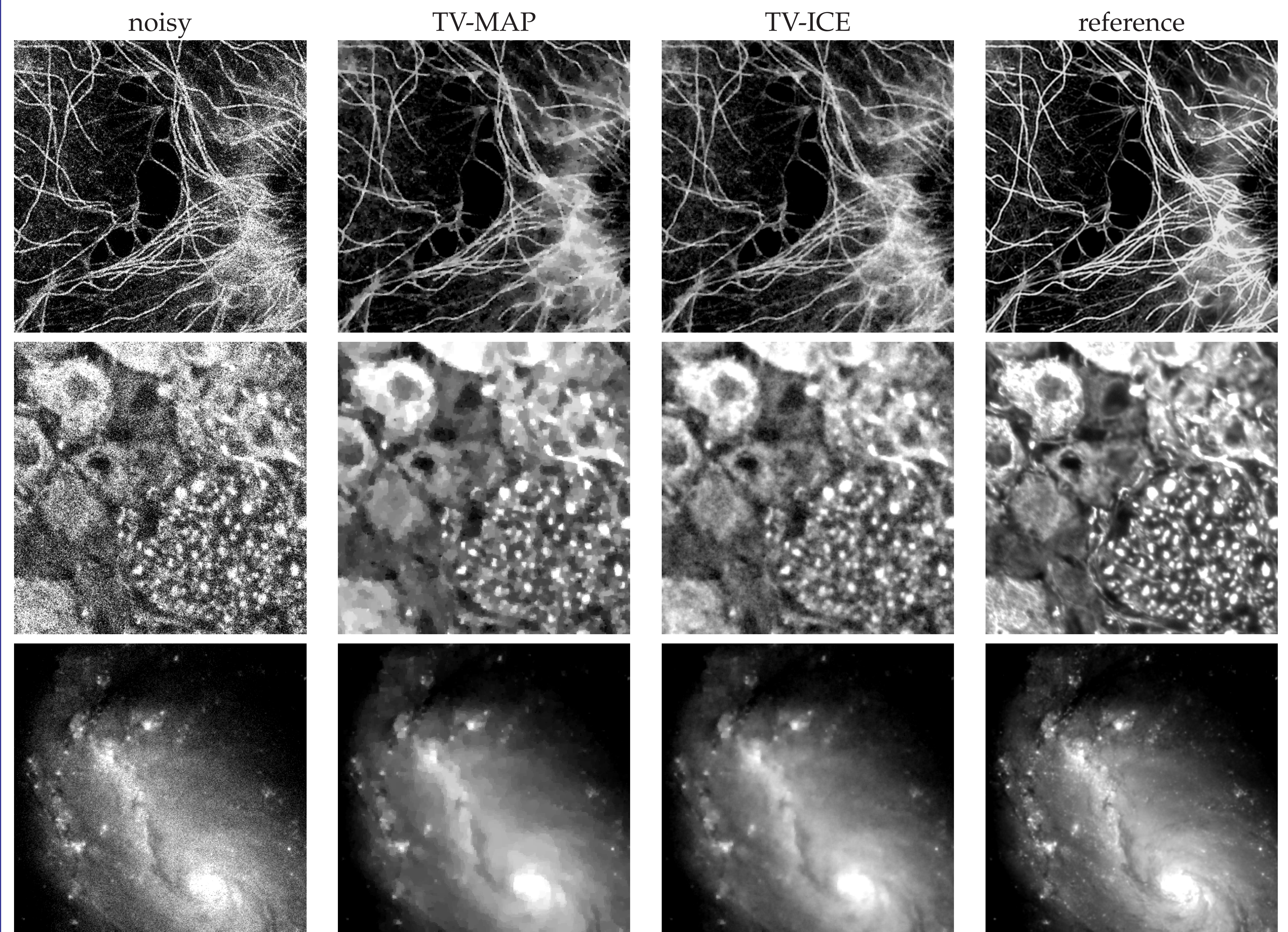
$$\gamma_\mu(p, x) = \int_0^x s^{p-1} e^{-\mu s} ds \quad (\forall \mu)$$

$$\Gamma_\mu(p, x) = \int_x^{+\infty} s^{p-1} e^{-\mu s} ds \quad (\mu > 0)$$

### Numerical computation: $\gamma_\mu$ or $\Gamma_\mu$ ?



### EXPERIMENTS



## Mathematical properties

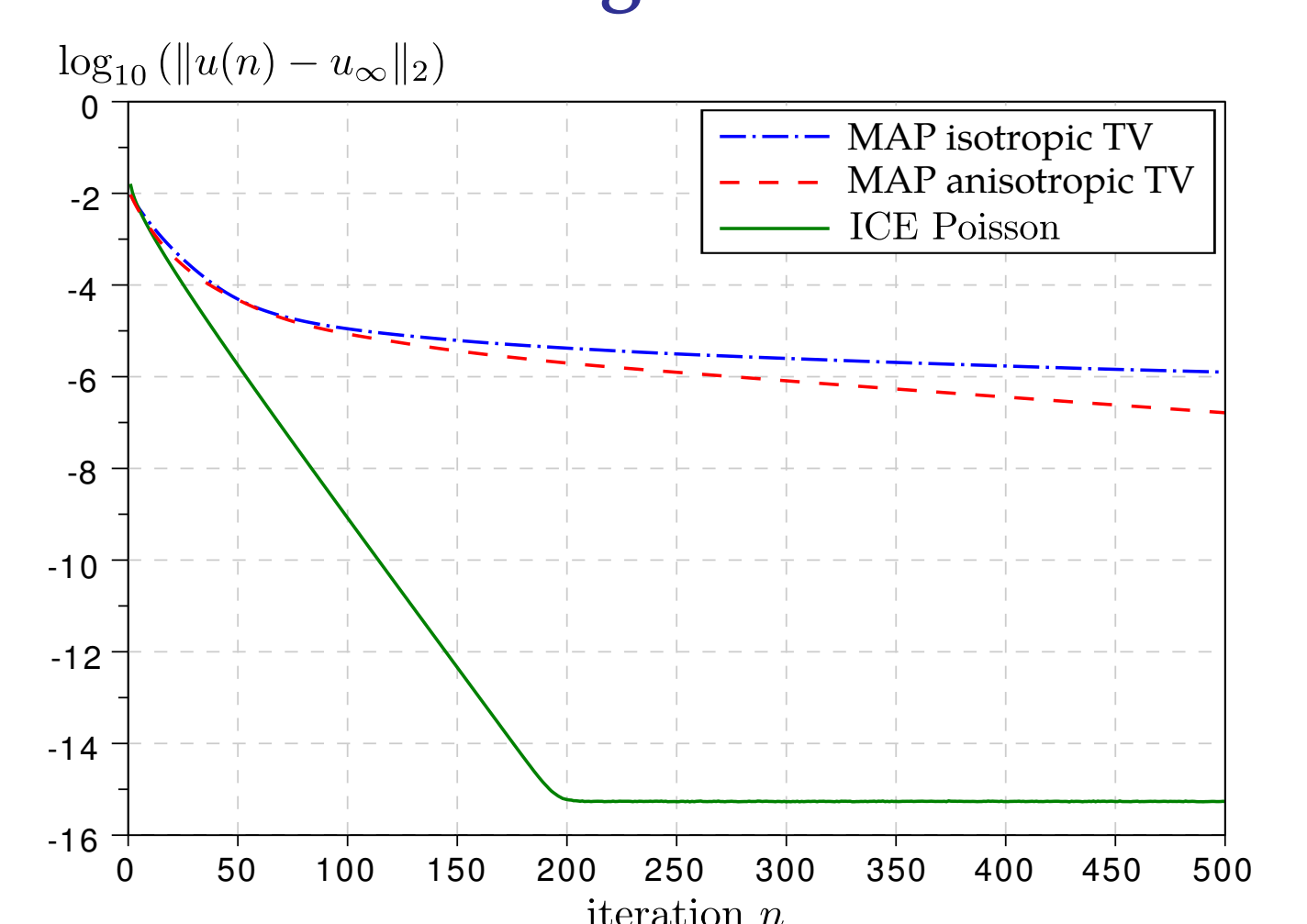
**Theorem 1 (linear convergence).** Given an image  $v \in \mathbb{R}^\Omega$ , the sequence of images  $(u^n)_{n \geq 0}$  defined by  $u^0 = 0$  and the TV-ICE recursion **converges linearly** to an image  $\hat{u}_{\text{ICE}}$ .

**Theorem 2 (no staircasing).** Let  $v : \Omega \rightarrow \mathbb{N}$  be a noisy image, and  $\hat{u}_{\text{ICE}}$  its denoised version. Let  $x$  and  $y$  be two pixels of  $\Omega$ , then if  $\hat{u}_{\text{ICE}}$  is constant over  $\mathcal{N}_x \cup \mathcal{N}_y \cup \{x, y\}$ , necessarily  $v(x) = v(y)$ .

More intuitively, Theorem 2 proves that Poisson **TV-ICE cannot produce large constant regions** that were not at least partially present in the initial data.

**Main perspective:** use the Poisson TV-ICE model in **more complex inverse problems** involving TV terms.

### Convergence rate



## Main references

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- Louchet, C., Moisan, L.: "Posterior Expectation of the Total Variation model: Properties and Experiments". SIAM J. Imaging Sci. (2013)
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